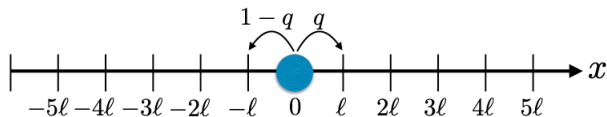
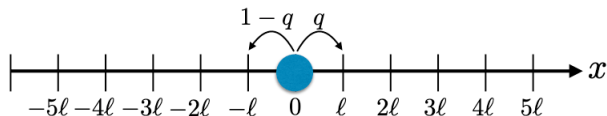


## A 1 dimensional random walk



A particle starts at  $x = 0$  and jumps left or right with probability  $q$ , left with probability  $1 - q$ .

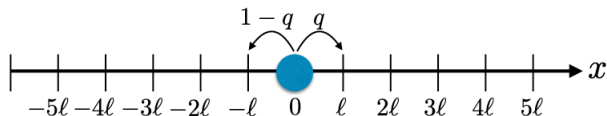
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What is the probability that we find the particle at position  $x$  after  $n$  jumps?

The particle is at position  $x$ , iff it has made  $k$  jumps to the right where  $x = k - (n - k)$ . The probability of having jumped  $k$  times to the right is a binomial:

$$P(k, n) = \binom{n}{k} q^k (1 - q)^{n-k}$$

## Deriving diffusion

According to the central limit theorem, the Binomial distribution  $P(k, n)$  approaches the Gaussian with mean  $\mu = n(2q - 1)$  and variance  $\sigma^2 = 4nq(1 - q)$ <sup>1</sup>:

$$p(x, n) \approx \mathcal{N}\left(n(2q - 1), 4nq(1 - q)\right).$$

---

<sup>1</sup>For  $q = \frac{1}{2}$ , it's  $\mathcal{N}(0, n)$ .

## Scaling the master equation

The master equation gives a recursive relation between the probability distributions at time  $n$ :

$$P(x, n + 1) = qP(x - 1, n) + (1 - q)P(x + 1, n)$$

$$P(x, 0) = \delta_0(x).$$

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What if we speed up the process and make the step size shrink? The correct scaling<sup>2</sup> is to let  $\Delta x = \sqrt{\Delta t}$ .

Derivation: write  $p + \delta t$

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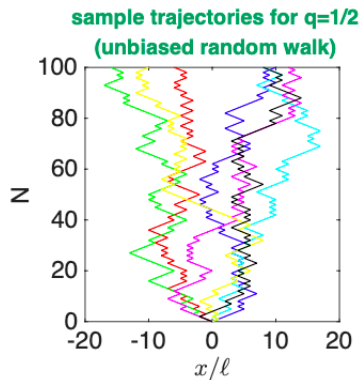
<sup>2</sup>Since the variance of the position grows at rate  $\sqrt{\Delta t}$  

## With drift

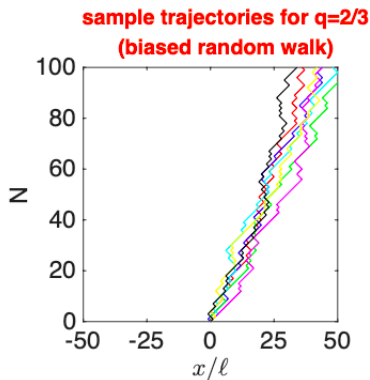
What if  $q \neq \frac{1}{2}$ ? Then, there is some drift involved:

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4



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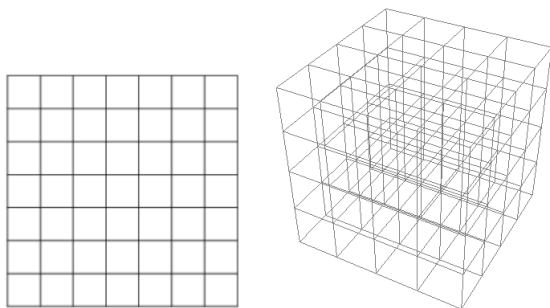
The Fokker-Planck equation is like the heat equation, but with drift:

$$p_t = -vp_x + \Delta p$$

Derive this by Taylor-expanding the master equation,

## Diffusion in $n$ dimensions

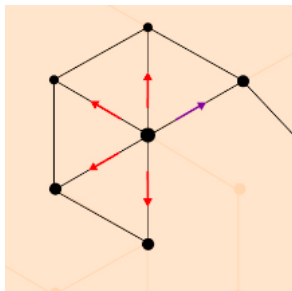
All of this can be generalized to  $n$  dimensions, of course.



- ▶ Particle moves in each direction with probability  $\frac{1}{2^n}$ .
- ▶ Shizuo Kakutani: “A drunk man will find his way home, but a drunk bird may get lost forever.”

# Graph random walk

What about a random walk on the nodes of a graph?



The movement of the random walker has to do with the view of the adjacency matrix as a linear operator!

## Graph random walk

The normalized adjacency matrix is

$$W(i,j) = \frac{1}{\deg(i)} A(i,j)$$

E.g.

$$A = \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} \Rightarrow W = \begin{pmatrix} 0 & 1 & 0 \\ \frac{1}{2} & 0 & \frac{1}{2} \\ 0 & 1 & 0 \end{pmatrix}$$

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Row  $i$  can be thought of as a discrete probability distribution of the next position given that we start at node  $i$ .

# Graph random walk

How does the normalized adjacency matrix  $W$  behave as a linear operator?

$$W : \mathbb{R}^n \rightarrow \mathbb{R}^n$$

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Let  $v = e_k$  for some basis vector  $e_k$  of  $\mathbb{R}^n$ . We can think of this as an initial state: a particle at node  $k$ .

Then  $Wv$  represents the probability distribution after one walker step.  $W^t$  represents the probability distribution after  $t$  steps.  $W$  is a Markov transition matrix.

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$v$  can also be any normalized vector.

# Graph random walk

We can apply the theory of stochastic processes and Markov chains to the graph random walk problem.

- ▶ Existence of stationary distributions
- ▶ Convergence to stationary distributions
- ▶ Eigenvalues and eigenvectors of  $W$ . Property: its largest eigenvalue is always  $\lambda = 1$  and all the rest have  $|\lambda_i| < 1$ .

## Application: PageRank

Let the graph be the (directed) graph of hyperlinks on the web. Each vertex is a webpage, with edges pointing towards the websites it links.

---

<sup>3</sup>The algorithm actually has some stochasticity, where we click a link with probability  $\alpha$  and go to a completely new webpage with probability  $1 - \alpha$ .

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The stationary distribution gives us a ranking of the most influential/popular pages on the internet.

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# Diffusion equation

$$u_t = \Delta u.$$

---

<sup>4</sup>I tried to derive the graph heat equation from the random walk/taylor expansion method, but it was unclear what scaling to use.

# Diffusion equation

$$u_t = \Delta u.$$

$$\Delta u = \operatorname{div} \nabla u(x) = \partial_{x_1} u + \partial_{x_2} u + \cdots + \partial_{x_n} u$$

What does the graph version look like?<sup>4</sup>

---

<sup>4</sup>I tried to derive the graph heat equation from the random walk/taylor expansion method, but it was unclear what scaling to use.

## Graph diffusion

The directional derivatives are replaced by derivatives in the direction of each neighboring node; the graph analogue of  $\operatorname{div} \nabla u$  is then

$$\begin{aligned}\Delta u(i) &= \sum_{j \in N(i)} u_j - u_i \\ &= \sum_{j=1}^n A_{ij}(u_j - u_i) \\ &= \sum_{j=1}^n A_{ij} u_j - u_i \operatorname{deg}(i) \\ &= (A - D)u\end{aligned}$$

where  $D = \operatorname{diag}(\operatorname{deg}(i))$ .

$A - D =: L$  is called the graph Laplacian.

# Graph diffusion

What do we do with this?

- ▶  $u_t = Lu$  is a linear system of ODEs, so is solvable.
- ▶ Since  $L$  is symmetric: real eigenvalues, and orthogonal eigenvectors
- ▶  $L$  is singular with eigenvalue 0 and corresponding eigenvector  $\mathbf{1}$ .

$$\vec{u}(t) = \sum_{i=1}^n a_i(0) e^{-c\lambda_i t} \vec{v}_i$$

Spectral Solution

## Graph diffusion: uses

The eigenvalues of the graph Laplacian tell us a lot about the structure of the graph

- ▶  $\lambda_1 = 0$  always. The number of eigenvectors corresponding to  $\lambda_1$  is the number of connected components of the graph.
- ▶ All eigenvalues must be nonnegative<sup>5</sup>

---

<sup>5</sup>If there was a negative e-value, then the corresponding term in the spectral solution would go to  $\infty$  which is not possible



diffusion from random walk



diffusion on graphs